

Lecture 10:Applications to eigenfunctions of the Laplacian:

$\Omega \subseteq \mathbb{R}^2$ nice planar domain with piecewise smooth boundary

$$\partial\Omega. \quad \Delta = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}$$

acting on functions vanishing near $\partial\Omega$.

Self adjoint operator Δ :

$$\int_{\Omega} f \Delta g = \int_{\Omega} \Delta f g \quad \text{"Dirichlet Laplacian"}$$

Fact:

$L^2(\Omega)$ has ONB of eigenfunctions of Δ .

Example: $\Delta = \frac{d^2}{dx^2}$ on $L^2(\mathbb{R})$. Self adjoint $\int f g'' = \int f'' g$.

No eigenfunctions in $L^2(\mathbb{R})$.

$$\Delta f = f'' = -k^2 f + f = a_+ e^{ikx} + a_- e^{-ikx} \notin L^2$$

$k \in \mathbb{R}, -k^2 < 0$

$$L^2(\mathbb{R}) \ni f(x) = \int \hat{f} e^{ikx} dk$$

Example 1:

dim 1, $\Omega = [0, A]$, $f_n(x) = \sqrt{\frac{2}{A}} \sin\left(\pi n \frac{x}{A}\right)$, $n \geq 1$

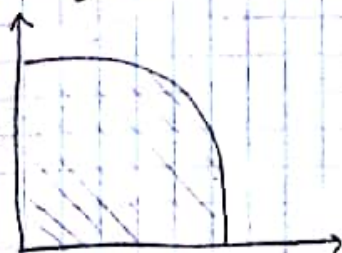
$f_n''(x) = -\left(\frac{\pi n}{A}\right)^2 f_n(x)$, $f_n|_{\partial\Omega} = 0$. ONB of $\{f \in L^2\}$ vanish on $\partial\Omega$.

Weyl's Law (1911):

For such $\Omega \subseteq \mathbb{R}^2$ $\#\{E_n \leq x\} \sim \frac{\text{area}(\Omega)}{4\pi} x$, $x \rightarrow \infty$.

e.g. $\Omega = A \begin{array}{|c|} \hline B \\ \hline \end{array}$ is $f_{mn} = \sqrt{\frac{2}{AB}} \sin\left(\pi m \frac{x}{A}\right) \sin\left(\pi n \frac{y}{B}\right)$

eigenvalues $E_{mn} = \left(\frac{\pi m}{A}\right)^2 + \left(\frac{\pi n}{B}\right)^2$, $m, n \geq 1$.



$$\left(\frac{\pi x}{A}\right)^2 + \left(\frac{\pi y}{B}\right)^2 \leq E \leftarrow \text{count solutions}$$

"hearing the area of Ω "

Thm. (Kac, McKean & Singer)

If Ω has smooth boundary then

$$\underbrace{\sum_{n=1}^{\infty} e^{-tE_n}}_{\text{trace of the heat kernel}} \underset{t \rightarrow 0}{\sim} c_1 \frac{\text{area}(\Omega)}{\sqrt{t}} - c_2 \frac{\text{length}(\partial\Omega)}{\sqrt{t}} + \frac{1-h}{6} + o(1)$$

$h = \# \text{holes}$

In general: M is a d -dim compact smooth Riemann manifold, then $-\Delta_g f_n = \lambda_n f_n$.

$$\sum_{n \geq 0} e^{-tE_n} = \frac{1}{(4\pi t)^{d/2}} \left\{ \text{Vol}(M) + a_1 t + a_2 t^2 + \dots \right\}$$

dim $d = 2$

$$\sum_{n \geq 0} e^{-tE_n} \underset{t \rightarrow 0}{\sim} \frac{1}{4\pi t} \left\{ \text{area}(M) + \frac{1-g(M)}{12} + O(t^2) \right\}$$

$g(M)$ = genus of M .

Let's compute the small time asymptotics of

$$K(t) = \sum e^{-tE_n} \text{ for a rectangle } \Omega_{A,B} = \underbrace{A \times B}_B$$

Method: reduce to 1-dim:

Eigenvalues $E_{m,n} = \left(\frac{\pi m}{A}\right)^2 + \left(\frac{\pi n}{B}\right)^2$

$$K(t) = \sum_{m,n=1}^{\infty} e^{-tE_{m,n}} = \sum_{m=1}^{\infty} e^{-t(\pi m/A)^2} \sum_{n=1}^{\infty} e^{-t(\pi n/B)^2} = K_{\Omega_A}(t) K_{\Omega_B}(t)$$

$\Omega_A = [0, A]$ interval.

And we get $\sum \sum e^{-\tau(n^2+m^2)} = \left(\sum e^{-\tau n^2} \right) \left(\sum e^{-\tau m^2} \right)$

1-dim $K_{[0,A]}(t) = \sum_{n=1}^{\infty} e^{-t(\pi n/A)^2} = \sum_{n=1}^{\infty} e^{-\pi n^2 \tau} - 1 \quad \ominus \quad \left(\tau = \frac{\pi}{A^2} t \right)$

When $\Theta(\tau) = \sum_{n \in \mathbb{Z}} e^{-\pi n^2 \tau}$ Theta function

$$\ominus \frac{\Theta(\tau) - 1}{2}$$

Note: as $\tau \rightarrow \infty$ $\Theta(\tau) = 1 + O(e^{-\pi \tau})$

$$\sum_{n=1}^{\infty} e^{-\tau n^2} \leq \sum_{k=1}^{\infty} e^{-\tau k} = \frac{e^{-\tau}}{1 - e^{-\tau}}$$

Thm.:

$$\Theta\left(\frac{1}{t}\right) = \sqrt{t} \Theta(t)$$

Cor.:

$$\Theta(t) \sim \frac{1}{\sqrt{t}} + o(e^{-\pi/t}) \quad , \quad t \gg 0$$

$$\Theta(t) = \frac{1}{\sqrt{t}} \Theta\left(\frac{1}{t}\right) = \frac{1}{\sqrt{t}} (1 + o(e^{-\pi/t}))$$

$\begin{matrix} t \gg 0 \\ 1/t \gg 0 \end{matrix}$

Cor.:

The 1-dim heat kernel

$$K_{[0,1]}(t) = \frac{\Theta(t) - 1}{2} = -\frac{1}{2} + \frac{1}{2} \left\{ \frac{1}{\sqrt{\pi t}} + \text{small} \right\}$$

$$K_{[0,1]}(t) = \frac{1}{2\sqrt{\pi}} \cdot \frac{1}{\sqrt{t}} - \frac{1}{2} + \text{small}$$

For rectangle:

$$\begin{aligned} K_{\Omega, \Omega}(t) &= K_{\Omega_A}(t) \cdot K_{\Omega_B}(t) = \left(\frac{1}{2\sqrt{\pi}} \cdot \frac{1}{\sqrt{t}} - \frac{1}{2} + \dots \right) \left(\frac{1}{2\sqrt{\pi}} \cdot \frac{1}{\sqrt{t}} - \frac{1}{2} + \dots \right) = \\ &= \frac{AB}{4\pi} \cdot \frac{1}{t} - \frac{A+B}{4\sqrt{\pi}} \cdot \frac{1}{\sqrt{t}} + \frac{1}{4} + o(1) = \\ &= \frac{\text{Area}(\Omega)}{4\pi} \cdot \frac{1}{t} - \frac{\text{length}(\partial\Omega)}{8\sqrt{\pi}} \cdot \frac{1}{\sqrt{t}} + \frac{1}{4} + o(1) \end{aligned}$$

Compare: In general smooth $\Omega \subseteq \mathbb{R}^2$,

$$K_{\Omega}(t) = \frac{\text{area}(\Omega)}{4\pi} \cdot \frac{1}{t} - \frac{1}{8\sqrt{\pi}} \frac{\text{length}}{\sqrt{t}} + o(1)$$

Notice: A rectangle is not smooth.

Lets prove that $\Theta(1/t) = \sqrt{t} \Theta(t)$

proof:

$$\Theta(t) = \sum_{n \in \mathbb{Z}} e^{-\pi t n^2} \quad . \quad \text{Let } g(x) = e^{-\pi x^2}$$

$t > 0$, abs. convergent. $g_t(x) = g(\sqrt{t}x) = e^{-\pi t x^2}$

$$\Theta(t) = \sum_{n \in \mathbb{Z}} g_t(n) = \sum_{m \in \mathbb{Z}} \hat{g}_t(m)$$

Poisson Summation

Recall $\forall y, g_\lambda(y) = g(\lambda y)$ then $\hat{g}_\lambda(y) = \frac{1}{\lambda} \hat{g}(\frac{y}{\lambda})$.

Also for $g(x) = e^{-\pi x^2}$, $\hat{g} = g \Rightarrow$

$$\hat{g}_{\sqrt{c}}(x) = \frac{1}{\sqrt{c}} \hat{g}\left(\frac{x}{\sqrt{c}}\right) = \frac{1}{\sqrt{c}} g\left(\frac{x}{\sqrt{c}}\right) = \frac{1}{\sqrt{c}} e^{-\pi x^2/c}$$

$$\Theta(\tau) = \sum_{m \in \mathbb{Z}} \hat{g}_{\sqrt{c}}(m) = \sum_{m \in \mathbb{Z}} \frac{1}{\sqrt{c}} e^{-\pi m^2/c} = \frac{1}{\sqrt{c}} \Theta(1/\tau)$$

Ex.: $\Omega: \pi \left\{ \begin{array}{c} \square \\ \hline \square \end{array} \right\}$ $\sin_{m,n}(x,y) = \sin(mx) \sin(ny)$

$$Z(f_{m,n}) = \begin{array}{c} \square \\ \square \\ \square \\ \square \end{array}$$

However, if $A=B$ then the eigen space have $\dim > 1$.

eigen value $E_{m,n} = \left(\frac{\pi m}{a}\right)^2 + \left(\frac{\pi n}{b}\right)^2 = m^2 + n^2$
 $a=b=\pi$

$$\dim \text{eigenspace} = \# \{ (m,n) \in \mathbb{N}^2 \mid E = m^2 + n^2 \}$$

If $E = p \equiv 1 \pmod{4}$ prime then $r_1(p) = 8$.

Diameter of p eigenspace = 2.

$r_2(65) = 16$ so $\dim(65) = 4$ so we get complicated

eigenfns. $f_E(x,y) = \sum_{\substack{m^2+n^2=E \\ m,n \geq 1}} a_{m,n} \sin(mx) \sin(ny)$

General theory:

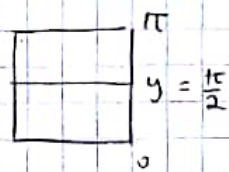
$$C_1 \sqrt{E} \leq \text{length}(Z(f_E)) \leq C_2 \sqrt{E}$$

$\exists C_1(\Omega) < C_2(\Omega)$ s.t. $\forall$ eigenfunction $-\Delta f = E$. $Z(f)$ is smooth outside of finite many pts. & at singular points it looks like equiangular intersecting.

Persistent components:

Does the model set vary with $\varepsilon \rightarrow \infty$?

Example: the line $y = \frac{\pi}{2}$ is on the model set S .



$$f_{m,n}(x,y) = \sin(mx)\sin(2ny)$$

$$\forall m,n \geq 1, E = m^2 + (2n)^2$$

Def.:

Let $\Sigma \subseteq \Omega_{\pi,\pi} = \square$ be a fixed curve. We say Σ is persistent if $\exists E_n \rightarrow \infty$ & eigenfunctions f_n with eigenvalues E_n s.t. $f_n|_{\Sigma} \equiv 0 \quad \forall n$.

Example: $\Sigma = \{y = \pi/2\}$

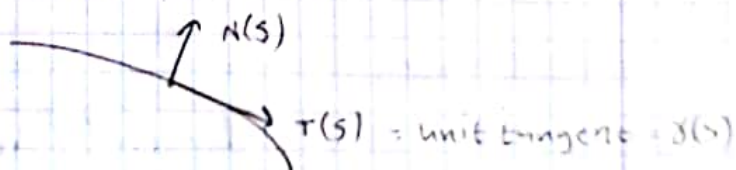
This is essentially the only example!

Thm.:

Assume $\Sigma \subset \square$ is a real-analytic curve with nowhere vanishing curvature. Then Σ is not persistent, i.e. $\exists E_{\Sigma} > 0$ s.t. if $E \geq E_{\Sigma}$ & f_E is an eigenfunction with eigenvalue E then $f_E|_{\Sigma} \neq 0$.

For $\gamma: [0,L] \rightarrow \Sigma$ arclength parameterization so

$$\|\dot{\gamma}(s)\| = 1$$



Def 1:

Curvature $K(s)$ of the pt. $\gamma(s)$ is $\dot{T}(s) = K(s)N(s)$.

Def 2:

Let $\alpha(s)$ be the angle between $T(s)$ & the x axis. Then

$$K(s) := \left| \frac{d\alpha}{ds} \right|$$

F-S eqs.:

$$\dot{T} = \kappa N, \quad \dot{N} = -\kappa T$$

We assume Σ has nowhere zero curvature:

$$\kappa(s) \neq 0 \quad \forall s \in [0, 1] \quad 0 < \kappa_{\min} \leq \kappa(s) \leq \kappa_{\max}$$

Lemma:

Let $\Sigma \subset \mathbb{R}^2$ be a real analytic curve with nowhere vanishing curvature. Let $\gamma: [0, 2] \rightarrow \Sigma$ be an arc length parameterization. $I_{\Sigma}(\xi) = \int_0^2 e^{i \langle \xi, \gamma(t) \rangle} dt$. Then

$$\exists C_{\gamma} > 0 \text{ s.t. } |I_{\Sigma}(\xi)| \leq C_{\gamma} \frac{1}{|\xi|^{1/2}}$$

Example: $\Sigma = \text{unit circle}$ $\gamma(s) = (\cos(s), \sin(s))$

$$I_{\Sigma}(\xi) = \int_0^{2\pi} e^{i \xi \cos(t)} dt$$

Proof of theorem:

$$\vec{x} = (x, y)$$

$$-\Delta f = E f \quad f = \sum_{\substack{m, n \in \mathbb{Z} \\ m, n \geq 1}} a_{m,n} \sin(mx) \sin(ny) = \sum_{\substack{|\xi|^2 = E \\ \xi \in \mathbb{Z}^2}} a(\xi) e^{i \langle \xi, \vec{x} \rangle}$$

choose $\xi_0 = (m_0, n_0)$ s.t. $|a(\xi_0)| \geq |a(\xi)| \quad \xi \in \mathbb{Z}^2 \quad \forall \xi$

After dividing by $a(\xi_0)$ we have

$$f = e^{i \langle \xi_0, \vec{x} \rangle} + \sum_{\substack{|\xi|^2 = E \\ \xi \neq \xi_0}} a(\xi) e^{i \langle \xi, \vec{x} \rangle} \quad \text{notice } |a(\xi)| \leq 1$$

$$\text{Let } J = \int_{\Sigma} e_{\xi_0} \cdot f = \int_0^2 e^{-i \langle \xi_0, \gamma(t) \rangle} f(\gamma(t)) dt$$

$f|_{\Sigma} \equiv 0 \Rightarrow J = 0$. On the other hand

$$\begin{aligned} J &= \int_0^2 e^{-i \langle \xi_0, \gamma(t) \rangle} \left(e^{i \langle \xi_0, \gamma(t) \rangle} + \sum_{\substack{|\xi|^2 = E \\ \xi \neq \xi_0}} a(\xi) e^{i \langle \xi, \gamma(t) \rangle} \right) dt \\ &= 2 + \sum_{\substack{|\xi|^2 = E \\ \xi \neq \xi_0}} a(\xi) \underbrace{\int_0^2 e^{i \langle \xi - \xi_0, \gamma(t) \rangle} dt}_{I(\xi - \xi_0)} \end{aligned}$$

By vdc lemma $I(\xi - \xi_0) \ll \frac{1}{|\xi - \xi_0|^{1/2}}$

By Jarnik's Thm, there is at most one $\xi_1 \neq \xi_0$ s.t.
 $|\xi_1 - \xi_0| < (\sqrt{\epsilon})^{1/3} = \epsilon^{1/6}$.

The sum over $\xi \neq \xi_0, \xi_1$ is bounded by:

$$\ll \sum_{\substack{\xi \neq \xi_0, \xi_1 \\ |\xi|^2 = \epsilon}} |a(\xi)| \frac{1}{|\xi - \xi_0|^{1/2}} \ll \sum 1 \cdot \frac{1}{(\sqrt{\epsilon})^{1/3}} \leq \\ \leq \frac{\#\{\xi \mid \xi^2 = \epsilon\}}{\epsilon^{1/6}} \ll \frac{\epsilon^\epsilon}{\epsilon^{1/6}} \quad \forall \epsilon > 0$$

So we find $0 = J = 2 + a(\xi_1) \cdot I(\xi_1 - \xi_0) + O\left(\frac{1}{\epsilon^{t-\epsilon}}\right)$

If ξ_1 does not exist we get a contradiction for $\epsilon \gg 1$.

Remains to deal with $I(\xi_1 - \xi_0)$.

Trivial bound: $|I(\xi)| \leq 2$. $I(\xi) = \int_0^\infty e^{i\langle \xi, \delta(t) \rangle} dt$.

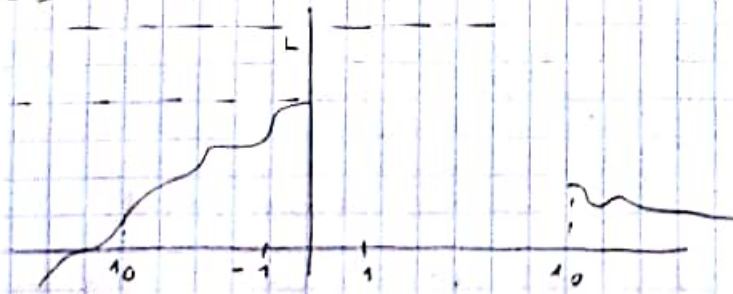
Know: $|\xi_1 - \xi_0| \geq 1$.

Want to avoid $|I(\xi_1 - \xi_0)| = 2$

Claim:

If $|V| \geq 1$ then $|I(V)| \leq cL$ (Then we win!)

It suffices to show $|I(V)| \neq 2$ for $V \neq 0$.



Because $I(V)$ is continuous, $|I(V)| \leq \frac{1}{2}$ for $|V| \geq 10$ &
 by compactness for $1 \leq |V| \leq 10$, $\max_{1 \leq |V| \leq 10} |I(V)| \leq 2$
 so is $\leq \frac{9}{10} \cdot 2$

Cauchy Schwarz:

$$\left| \int_0^L f(t) g(t) dt \right| \leq \sqrt{\int_0^L f^2} \sqrt{\int_0^L g^2}$$

& equality $\iff g = \lambda f, \lambda \in \mathbb{C}$.

Apply to $f(t) = 1, g(t) = e^{i \langle v, \gamma(t) \rangle}$ So

$$\int_0^L |f|^2 = \int_0^L |g|^2 = L$$

So $|I| = L \iff e^{i \langle v, \gamma(t) \rangle} \equiv \lambda \text{ constant} \iff$

$$\iff 0 = i \langle v, \dot{\gamma}(t) \rangle e^{i \langle v, \gamma(t) \rangle}$$

$$\frac{d}{dt} = 0 \iff \langle v, \dot{\gamma}(t) \rangle \equiv 0 \quad \forall t$$

So $\dot{\gamma}(t) \equiv u$ is fixed $\implies \gamma(t) = tu + u_0$ is a straight line