

W4 D 3

18/11/14

Random Spanning Trees

Reminder

$\langle \theta, \theta' \rangle = \sum_{e \in E} r_e \theta(e) \theta'(e) = \sum_{e \in E} r_e \theta(e) \theta'(e)$
 $\downarrow$
 $\ell^2(E, \{r_e\})$ with a basis $\chi^e = 1_{\vec{e}} - 1_{\vec{e}'}$,
 orthogonal but not necessarily of norm 1.

a star at v is $\sum_{e^+ = v} \chi^e$ STAR = span(star(v)).

Given a directed cycle $e_1, \dots, e_n$ in G
 the $\ell^2(E)$ element $\sum_{i=1}^n \chi^{e_i}$ is the cycle in $\ell^2(E)$. We define CYCLE = span(cycles).

In finite G we observed STAR + CYCLE = $\ell^2(E)$.

If $\theta: a \rightarrow \mathbb{Z}$ flow $\begin{bmatrix} d^* \theta(v) = 0 \quad \forall v \neq a, z \\ d^* \theta(a) = -d^* \theta(z) \end{bmatrix}$. We
 get a current flow by projecting it to STAR
 with the same strength

Reciprocity

For two distinct edges $e = e'$ (so $\vec{e} \neq \vec{e}' \neq \vec{e}$),

$$\langle P_{\text{STAR}} \chi^e, \chi^{e'} \rangle = \langle i^e, \chi^{e'} \rangle = i^e(e') r_{e'}.$$

i^e = unit current flow from e^- to e^+ .

But P_{STAR} is self adj., so it obviously
 is equal to $\langle \chi^e, P_{\text{STAR}} \chi^{e'} \rangle = i^{e'}(e) r_e$.

Denote $y(e, e') = i^e(e') r_{e'}$, then
 $y(e, e') r_e = y(e'; e) r_e$. y is called the
Transfer current Matrix.

Contraction of $e \in E$ is a new graph G/e obtained
 by removing e and identifying v_-, v_+ . $G \setminus e$ will mark
 G after deleting e . T_e = UST of G .

Note T_e given $e \in T_e$ is distributed like G/e
 " " " " " like $G \setminus e$

example $P(e, f \in T) = P(e \in T) P(f \in T | e \in T) = P(e \in T) P(f \in T / e)$

Cor. $P(e, f \in T) \leq P(e \in T) P(f \in T)$, by Kirchhoff's formula,
 & Rayleigh Monotonicity.

What does contraction do to $\ell^2(E)$?

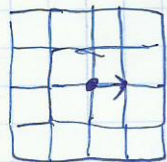
Let i_e be the UCF in G , $e^- \rightarrow e^+$. Let $F \subseteq E$ s.t. $e \notin F$, and won't close a loop ~~in~~ in $(G \setminus F)$. Write i_e for UCF $e^- \rightarrow e^+$ in G_{min}/F .

Let $Z = \text{span}\{i_f : f \in F\}$.

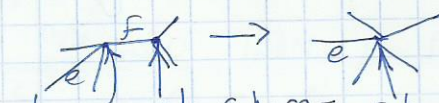
Claim: $i_e = P_{Z^\perp} i_e$. We'll check $P_Z i_e = 0$. [or, if we keep contracted edges has loops. $\hat{i} = P_Z i_e$.

Pf: since $i_e \in \text{STAR}$, $P_Z i_e \in \text{STAR}$, $P_Z i_e \in \text{STAR}$ so the cycle law holds. To verify $P_Z i_e = 0$, $\langle P_Z i_e, i_f \rangle = \langle i_e, i_f \rangle = 0$.

To see the new law of $P_Z i_e$ holds in G/F we write $P_Z i_e = i_e - \sum_{f \in F} \alpha_f i_f$. Note that a star ψ in G/F is a sum $\oplus$ of stars in G s.t. $\partial(f) = 0$ for $\forall f \in F$. i_e is orthogonal to any star ~~at~~ at $v \in e^\pm$. So $i_e \perp \psi$. Similarly $i_f \perp \psi$ where e, f don't share a vertex, and ψ is not a star at e^- or e^+ , $P_Z i_e \perp \psi$.



mass out of e^- is 1 in i_e , 0 in i_f for a star in $e^\pm$. What if e, f share a vertex?



star here + star here = star here and the mass out of e^- is 1 in i_e , 0 in i_f .

A different pt. $i_e = P_{Z^\perp} i_e$ [G/F has same edges as $G \setminus F$ turned to]. Let $\widehat{\text{STAR}}$ and $\widehat{\text{CYCLE}}$ be the loops resp. spaces in G/F . We'll find $\widehat{\text{CYCLE}} = \widehat{\text{CYCLE}} + \text{span}\{i_f : f \in F\}$, so $\widehat{\text{CYCLE}} \subseteq \widehat{\text{CYCLE}}$ and hence $\widehat{\text{STAR}} \subseteq \widehat{\text{STAR}}$.

STAR	CYCLE
STAR	CYCLE

 $\widehat{\text{STAR}} \perp \widehat{\text{CYCLE}} \Rightarrow P_{\widehat{\text{STAR}}} \widehat{\text{CYCLE}} = \widehat{\text{STAR}} \cap \widehat{\text{CYCLE}}$, so $P_{\widehat{\text{STAR}}} \{i_f : f \in F\} = Z = \text{span}\{i_f : f \in F\}$ and $\widehat{\text{STAR}} = \widehat{\text{STAR}} \oplus Z$.

So $\ell^2(E) = \text{STAR} \oplus \text{CYCLE} \oplus \mathbb{Z}$. Now, $\hat{i}^e = P_{\text{STAR}} X^e = P_{\text{STAR}} P_{\mathbb{Z}} X^e$
 $\mathbb{Z}^+ \underbrace{P_{\text{STAR}}}_{i^e} X^e$

The transfer-current thm.

$$P(e_1, \dots, e_k \in T) = \det \{Y(e_i, e_j)\}_{1 \leq i, j \leq k}$$

pf. If $e_1, \dots, e_k$ contains a cycle then the LHS is obviously 0. Write $\sum_{j=1}^k a_j X^{e_j}$ ^{not all 0} $a_j \in \{-1, 0, 1\}$ for the cycle. $\forall m=1, \dots, k$ $\sum_{j=1}^k a_j e_j \cdot i^{e_m}(e_j) = 0$, so $\{Y(e_i, e_j)\}_{1 \leq i, j \leq k} \begin{pmatrix} e_i \cdot a_j \\ e_j \cdot a_j \end{pmatrix} = 0$.
 Proceed by induction $\left\{ \begin{array}{l} \text{cycle law for } i^{e_m} \end{array} \right.$

$k=1$ is Kirchhoff's formula. put $Y_m = \{Y(e_i, e_j)\}_{1 \leq i, j \leq m}$
 for $m=1, \dots, k$. $P(e_1, \dots, e_k \in T) = P(e_1, \dots, e_{k-1} \in T) P(e_k \in T | e_1, \dots, e_{k-1})$
 $= \det Y_{k-1} \cdot \hat{i}^{e_k}(e_k) \stackrel{?}{=} \det Y_k$. We know $\hat{i}^{e_k} = P_{\mathbb{Z}^+} i^{e_k}$, where
 $\mathbb{Z} = \text{span}\{i^{e_1}, \dots, i^{e_{k-1}}\}$. So $\hat{i}^{e_k} = i^{e_k} - \sum_{j=1}^{k-1} a_j i^{e_j}$.
 $Y_k = \begin{pmatrix} i^{e_1} \\ \vdots \\ i^{e_k} \end{pmatrix} \xrightarrow{i^{e_k} - \sum_{j=1}^{k-1} a_j i^{e_j}} \begin{pmatrix} i^{e_1} \\ \vdots \\ i^{e_k} - \sum_{j=1}^{k-1} a_j i^{e_j} \end{pmatrix} \rightarrow \begin{pmatrix} i^{e_1} \\ \vdots \\ i^{e_{k-1}} \end{pmatrix} \begin{pmatrix} i^{e_k} \\ \vdots \\ i^{e_k} \end{pmatrix}$. But

$$\hat{i}^{e_k} = (0, 0, 0, \dots, \hat{i}^{e_k}(e_k)) \text{ so } \det Y_k =$$

$$\det \begin{pmatrix} i^{e_1} \\ \vdots \\ i^{e_{k-1}} \\ 0 \dots 0 \hat{i}^{e_k}(e_k) \end{pmatrix} = \det Y_{k-1} \cdot \hat{i}^{e_k}(e_k), \text{ as we wanted to prove.}$$

We saw $P(f \in T | e \in T) \leq P(f \in T)$, for $e \neq f$. We can expect a ~~similar~~ similar result for any increasing event.